# Bounds for blow-up time in a semilinear parabolic problem with variable exponents 

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#### Abstract

This report deals with a blow-up of the solutions to a class of semilinear parabolic equations with variable exponents nonlinearities. Under some appropriate assumptions on the given data, a more general lower bound for a blow-up time is obtained if the solutions blow up. This result extends the recent results given by Baghaei Khadijeh et al. [8], which ensures the lower bounds for the blow-up time of solutions with initial data $\varphi(0)=\int_{\Omega} u_{0}{ }^{k} d x, k=$ constant.


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## 1. Introduction

In this paper, we are concerned with the following semilinear parabolic equation

$$
\left\{\begin{array}{c}
u_{t}-\Delta u=u^{p(x)}, \quad x \in \Omega, t>0  \tag{1.1}\\
u=0 \text { on } \Gamma, \quad t \geq 0 \\
u(x, 0)=u_{0}(x) \geq 0, x \in \Omega
\end{array}\right.
$$

where $\Omega$ be a bounded domain in $\mathbb{R}^{n}$, with a smooth boundary $\Gamma=\partial \Omega, T \in(0,+\infty]$, and the initial value $u_{0} \in H_{0}^{1}(\Omega)$, the exponent $p($.$) is given measurable function on$ $\bar{\Omega}$ such that:

$$
\begin{equation*}
1<p_{1}=\underset{x \in \Omega}{e s s \inf } p(x) \leq p(x) \leq p_{2}=\underset{x \in \Omega}{e s s} \sup p(x)<\infty \tag{1.2}
\end{equation*}
$$

and satisfy the following Zhikov-Fan uniform local continuity condition:

$$
\begin{equation*}
|p(x)-p(y)| \leq \frac{M}{|\log | x-y| |}, \text { for all } x, y \text { in } \Omega \text { with }|x-y|<\frac{1}{2}, M>0 \tag{1.3}
\end{equation*}
$$

The problem (1.1) arises from many important mathematical models in engineering and physical sciences. For example, nuclear science, chemical reactions, heat transfer,
population dynamics, biological sciences, etc., and have interested a great deal of attention in the research, see $[4,7,9]$ and the references therein. For problem (1.1), Hua Wang et al. [10] established a blow-up result with positive initial energy under some suitable assumptions on the parameters $p($.$) and u_{0}$. In [9], the authors proved that there are non-negative solutions with a blow-up in finite time if and only if $p_{2}>1$. The authors in [11] obtained the solution of problem (1.1) blows up in finite time when the initial energy is positive. The following problem was considered by R . Abita in [3]

$$
u_{t}-\Delta u_{t}-\Delta u=u^{p(x)}, \quad x \in \Omega, t>0
$$

The author proved that the nonnegative classical solutions blow-up in finite time with arbitrary positive initial energy and suitable large initial values. Also, he employed a differential inequality technique to obtain an upper bound for blow-up time if $p($. and the initial value satisfies some conditions. In [8], the authors based exactly on the idea on the one in [6], derived the lower bounds for the time of blow-up, if the solutions blow-up. In order to declare the main results of this paper, we need to add the following energy functional corresponding to the problem (1.1) (see [2])

$$
\begin{equation*}
E(t)=\frac{1}{2} \int_{\Omega}|\nabla u(t)|^{2} d x-\int_{\Omega} \frac{1}{p(x)+1} u(x, t)^{p(x)+1} d x . \tag{1.4}
\end{equation*}
$$

## 2. Lower bounds of the blow-up time

In this section, we investigate the lower bound for the blow-up time $T$ in some suitable measure. The idea of the proof of the following theorem is inspired by on the one in [6]. For this goal, we start by the following lemma concerning the energy of the solution.

Lemma 2.1. Let $u(x, t)$ be a weak solution of (1.1), then $E(t)$ is a nonincreasing function on $[0, T]$, that is

$$
\begin{equation*}
\frac{d E(t)}{d t}=-\int_{\Omega} u_{t}^{2}(x, t) d x \leq 0 \tag{2.1}
\end{equation*}
$$

and the inequality $E(t) \leq E(0)$ is satisfied.
We consider the following partition of $\Omega$,

$$
\Omega^{-}=\{x \in \Omega|1>(k(x)-1) \ln | u \mid\}, \quad \Omega^{+}=\{x \in \Omega|1 \leq(k(x)-1) \ln | u \mid\}, \forall t>0
$$ where each $\Omega^{ \pm}$depends on $t$, and setting

$$
\widetilde{E}(0)=\frac{1}{2} \int_{\Omega^{-}}\left|\nabla u_{0}\right|^{2} d x-\int_{\Omega^{-}} \frac{1}{p(x)+1} u_{0}^{p(x)+1} d x .
$$

Now, we are in a position to affirm our principal theorem results.
Theorem 2.2. Assume $u_{0} \in L^{k(.)}(\Omega)$, and the nonnegative weak solution $u(x, t)$ of problem (1.1) blows up in finite time $T$, then $T$ has a lower bound by:

$$
\begin{equation*}
\int_{\varphi(0)}^{+\infty} \frac{d \gamma}{C_{1}+C_{2} \gamma^{\frac{3 n-6}{3 n-8}}}, \tag{2.2}
\end{equation*}
$$

where

$$
\begin{equation*}
\varphi(0)=\int_{\Omega} \frac{1}{k(x)(k(x)-1)} u_{0}^{k(x)} d x \tag{2.3}
\end{equation*}
$$

where $k($.$) is a measurable function on \bar{\Omega}$ such that

$$
\begin{align*}
\max \left(1,2(n-2)\left(p_{2}-1\right)\right)<k_{1} & =\underset{x \in \Omega}{e \operatorname{ess} \inf } k(x) \leq k(x) \leq k_{2} \\
& =\underset{x \in \Omega}{\operatorname{ess} \sup } k(x)<\infty \tag{2.4}
\end{align*}
$$

and

$$
\begin{equation*}
\sqrt{C_{k}}=\sup _{x \in \bar{\Omega}}|\nabla k(x)| \in L^{2}(\Omega), C_{k}>0 \tag{2.5}
\end{equation*}
$$

and $C_{i}(i=1,2)$ are positive constants will be described later.
Notation 2.3. We note that the presence of the variable-exponent nonlinearities in (2.6) below, makes analysis in the paper somewhat harder than that in the related ones, we will establish and give a precise estimate for the lifespan $T$ of the solution in this case. The method used here is the differential inequality technique. However, our argument is considerably different and it is more abbreviated.

Proof of Theorem (2.2). Set

$$
\begin{equation*}
\varphi(t)=\int_{\Omega} \frac{1}{k(x)(k(x)-1)} u(x, t)^{k(x)} d x . \tag{2.6}
\end{equation*}
$$

Multiplying the equation Eq. (1.1) by $u$ and integrating by parts, we see

$$
\begin{aligned}
& \varphi^{\prime}(t)= \int_{\Omega} \frac{1}{k(x)-1} u^{k(x)-1} u_{t} d x=\int_{\Omega} \frac{1}{k(x)-1} u^{k(x)-1}\left(\Delta u+u^{p(x)}\right) d x \\
&=\int_{\Omega} \frac{1}{k(x)-1} u^{k(x)-1} \Delta u d x+\int_{\Omega} \frac{1}{k(x)-1} u^{k(x)+p(x)-1} d x \\
&=- \int_{\Omega} \nabla\left(\frac{1}{k(x)-1} u^{k(x)-1}\right) \nabla u d x+\int_{\Omega} \frac{1}{k(x)-1}|u|^{k(x)+p(x)-1} d x
\end{aligned}
$$

where we have used the divergence theorem, the boundary condition on $u$.
It is straightforward to check that

$$
\nabla\left(\frac{1}{k(x)-1} u^{k(x)-1}\right)=u^{k(x)}|u|^{-2} \nabla u+\frac{\nabla k(x)}{k(x)-1} u^{k(x)-1}\left(\ln |u|-\frac{1}{k(x)-1}\right)
$$

then, we get

$$
\begin{equation*}
\varphi^{\prime}(t)=-\int_{\Omega} u^{k(x)}|u|^{-2}|\nabla u|^{2} d x+\int_{\Omega} \frac{1}{k(x)-1} u^{k(x)+p(x)-1} d x+\mathcal{Q} \tag{2.7}
\end{equation*}
$$

where

$$
\mathcal{Q}=\int_{\Omega} u^{k(x)-1}\left(\frac{1}{(k(x)-1)^{2}}-\frac{1}{(k(x)-1)} \ln |u|\right) \nabla k(x) . \nabla u d x
$$

Considering the following properties of the function $\mathcal{G}$,

$$
\begin{gathered}
\mathcal{G}(\lambda)=\frac{\lambda^{\gamma}}{\gamma^{2}}(1-\gamma \ln \lambda), \quad 0 \leq \lambda \leq e^{\frac{1}{\gamma}} \\
\mathcal{G}(0)=\mathcal{G}\left(e^{\frac{1}{\gamma}}\right)=0, \mathcal{G}^{\prime}(\lambda)=-\lambda^{\gamma-1} \ln \lambda, \quad \max _{0 \leq \lambda \leq e^{\frac{1}{\gamma}}} \mathcal{G}(\lambda)=\mathcal{G}(1)=\frac{1}{\gamma^{2}},
\end{gathered}
$$

and using the fact that

$$
\int_{\Omega^{-}}|\nabla u|^{2} d x \leq 2 \widetilde{E}(0)+2 \int_{\Omega^{-}} \frac{1}{p(x)+1} u(x, t)^{p(x)+1} d x, \quad(\text { by }(1.4) \text { and }(2.1))
$$

applying the Hölder, Young inequalities and (2.5), $\mathcal{Q}$ is evaluated as follows:

$$
\begin{gather*}
\mathcal{Q}=\int_{\Omega} u^{k(x)-1}\left(\frac{1}{(k(x)-1)^{2}}-\frac{1}{k(x)-1} \ln |u|\right) \nabla k(x) . \nabla u d x \\
=\int_{\Omega \cap(1>(k(x)-1) \ln |u(x, t)|)} u^{k(x)-1}\left(\frac{1}{(k(x)-1)^{2}}-\frac{1}{k(x)-1} \ln |u|\right) \nabla k(x) . \nabla u d x \\
\int_{\Omega \cap(1 \leq(k(x)-1) \ln |u(x, t)|)} u^{k(x)-1}\left(\frac{1}{(k(x)-1)^{2}}-\frac{1}{k(x)-1} \ln |u|\right) \nabla k(x) . \nabla u d x \\
\leq \int_{\Omega^{-}} \frac{1}{(k(x)-1)^{2}}|u|^{k(x)-1}(1-(k(x)-1) \ln |u|)|\nabla u||\nabla k(x)| d x \\
\leq \int_{\Omega^{-}} \frac{1}{(k(x)-1)^{2}}|\nabla k(x)||\nabla u| d x \leq \frac{1}{2\left(k_{1}-1\right)^{2}}\left(C_{k}+\int_{\Omega^{-}}|\nabla u|^{2} d x\right) \\
\leq \frac{1}{2\left(k_{1}-1\right)^{2}}\left(C_{k}+2 E(0)+2 \int_{\Omega^{-}} \frac{1}{p(x)+1} u(x, t)^{p(x)+1} d x\right) \\
\leq \frac{1}{2\left(k_{1}-1\right)^{2}}\left(C_{k}+2 E(0)+\frac{2}{p_{1}+1} \max \left(\int_{\Omega^{-}}|u|^{p_{2}+1} d x, \int_{\Omega^{-}}|u|^{p_{1}+1} d x\right)\right) \\
\leq \frac{1}{\left(k_{1}-1\right)^{2}}\left(\frac{1}{2} C_{k}+E(0)+\frac{1}{p_{1}+1} e^{\frac{p_{2}+1}{k_{1}-1}}|\Omega|\right) . \tag{2.8}
\end{gather*}
$$

Because in $\Omega^{+}$, we have

$$
\int_{\Omega^{+}}|u|^{k(x)-1}\left(\frac{1}{(k(x)-1)^{2}}-\frac{1}{k(x)-1} \ln |u|\right)|\nabla k(x)| d x \leq 0
$$

while that of the first term in the right-hand side of (2.7) was estimated as follows

$$
-\int_{\Omega}|u|^{k(x)-2}|\nabla u|^{2} d x \leq-\min \left(\int_{\Omega}|u|^{k_{2}-2}|\nabla u|^{2} d x, \int_{\Omega}|u|^{k_{1}-2}|\nabla u|^{2} d x\right) .
$$

Using the fact

$$
\left|\nabla u^{\gamma}\right|=\gamma u^{\gamma-1}|\nabla u|
$$

to get

$$
\begin{equation*}
-\int_{\Omega}|u|^{k(x)-2}|\nabla u|^{2} d x \leq-\min \left(\frac{4}{\left(k_{2}\right)^{2}} \int_{\Omega}\left|\nabla u^{\frac{k_{2}}{2}}\right|^{2} d x, \frac{4}{\left(k_{1}\right)^{2}} \int_{\Omega}\left|\nabla u^{\frac{k_{1}}{2}}\right|^{2} d x\right) \tag{2.9}
\end{equation*}
$$

Plugging this estimate (2.8) and (2.9) into (2.7), we obtain

$$
\begin{align*}
\varphi^{\prime}(t) & \leq \min \left(\frac{-4}{\left(k_{2}\right)^{2}} \int_{\Omega}\left|\nabla u^{\frac{k_{2}}{2}}\right|^{2} d x, \frac{-4}{\left(k_{1}\right)^{2}} \int_{\Omega}\left|\nabla u^{\frac{k_{1}}{2}}\right|^{2} d x\right) \\
+ & \frac{1}{k_{1}-1} \int_{\Omega} u^{k(x)+p_{2}-1} d x+\frac{1}{k_{1}-1} \int_{\Omega} u^{k(x)+p_{1}-1} d x \\
& +\frac{1}{\left(k_{1}-1\right)^{2}}\left(\frac{1}{2} C_{k}+E(0)+\frac{1}{p_{1}+1} e^{\frac{p_{2}+1}{k_{1}-1}}|\Omega|\right) \tag{2.10}
\end{align*}
$$

By using (2.4), we can apply the Hölder and Young inequalities to get

$$
\begin{align*}
\int_{\Omega} u^{k(x)+p_{2}-1} d x & \leq \int_{\Omega} 1 \cdot \alpha_{1} d x+\int_{\Omega} \alpha_{2} \cdot u^{\frac{k(x)(2 n-3)}{2(n-2)}} d x  \tag{2.11}\\
& \leq\left(\sup \alpha_{1}\right)|\Omega|+\left(\sup \alpha_{2}\right)\left(\int_{\Omega} u^{\frac{k(x)(2 n-3)}{2(n-2)}} d x\right)
\end{align*}
$$

and

$$
\begin{align*}
\int_{\Omega} u^{k(x)+p_{1}-1} d x & \leq \int_{\Omega} 1 \cdot \alpha_{3} d x+\int_{\Omega} \alpha_{4} \cdot u^{\frac{k(x)(2 n-3)}{2(n-2)}} d x  \tag{2.12}\\
& \leq\left(\sup _{\Omega} \alpha_{3}\right)|\Omega|+\left(\sup _{\Omega}\right)\left(\int_{\Omega} u^{\frac{k(x)(2 n-3)}{2(n-2)}} d x\right)
\end{align*}
$$

where

$$
\begin{gathered}
\alpha_{1}=1-\frac{2(n-2)\left(k(x)+p_{2}-1\right)}{(2 n-3) k(x)}, \quad \alpha_{2}=\frac{2(n-2)\left(k(x)+p_{2}-1\right)}{(2 n-3) k(x)}, \\
\alpha_{3}=1-\frac{2(n-2)\left(k(x)+p_{1}-1\right)}{(2 n-3) k(x)}, \quad \alpha_{4}=\frac{2(n-2)\left(k(x)+p_{1}-1\right)}{(2 n-3) k(x)} ; \\
\text { observe that } \alpha_{2} \geq \alpha_{4} \text { and } \alpha_{1} \leq \alpha_{3} .
\end{gathered}
$$

Combining (2.11) and (2.12) with (2.10) give

$$
\begin{gather*}
\varphi^{\prime}(t) \leq \\
\frac{-1}{2} \frac{4}{\left(k_{2}\right)^{2}}\left(\int_{\Omega}\left|\nabla u^{\frac{k_{2}}{2}}\right|^{2} d x+\int_{\Omega}\left|\nabla u^{\frac{k_{1}}{2}}\right|^{2} d x\right) \\
+\frac{2}{k_{1}-1}\left(\sup _{\Omega} \alpha_{2}\right) \int_{\Omega} u^{\frac{k(x)(2 n-3)}{2(n-2)}} d x  \tag{2.13}\\
+\frac{1}{\left(k_{1}-1\right)^{2}}\left(\frac{1}{2} C_{k}+E(0)+\frac{1}{p_{1}+1} e^{\frac{p_{2}+1}{k_{1}-1}}|\Omega|\right)+\frac{|\Omega|}{k_{1}-1} \sup _{\Omega}\left(\alpha_{3}+\alpha_{1}\right)
\end{gather*}
$$

We now make use of Schwarz's inequality to the second term on the right-hand side of (2.13) as follows

$$
\begin{gather*}
\int_{\Omega} u^{\frac{k(x)(2 n-3)}{2(n-2)}} d x \leq\left(\int_{\Omega} u^{k(x)} d x\right)^{\frac{1}{2}}\left(\int_{\Omega} u^{\frac{k(x)(n-1)}{n-2}} d x\right)^{\frac{1}{2}}  \tag{2.14}\\
\quad \leq\left(\int_{\Omega} u^{k(x)} d x\right)^{\frac{3}{4}}\left(\int_{\Omega}\left(u^{\frac{k(x)}{2}}\right)^{\frac{2 n}{n-2}} d x\right)^{\frac{1}{4}}
\end{gather*}
$$

Next, by using the Sobolev inequality (see [5]), for $n \geq 3$, we get

$$
\begin{align*}
\left\|u^{\frac{k(x)}{2}}\right\|_{\frac{2 n}{n-2}}^{\frac{n}{2(n-2)}} & \leq B^{\frac{n}{2(n-2)}} \max \left(\left\|\nabla u^{\frac{k_{2}}{2}}\right\|_{2}^{\frac{n}{2(n-2)}},\left\|\nabla u^{\frac{k_{1}}{2}}\right\|_{2}^{\frac{n}{2(n-2)}}\right)  \tag{2.15}\\
& \leq B^{\frac{n}{2(n-2)}}\left(\left\|\nabla u^{\frac{k_{2}}{2}}\right\|_{2}^{\frac{n}{2(n-2)}}+\left\|\nabla u^{\frac{k_{1}}{2}}\right\|_{2}^{\frac{n}{2(n-2)}}\right)
\end{align*}
$$

where $B$ is the best constant in the Sobolev inequality.
By inserting the last inequality in (2.14) and (2.15), we have

$$
\begin{gathered}
\int_{\Omega} u^{\frac{k(x)(2 n-3)}{2(n-2)}} d x \leq \\
\leq B^{\frac{n}{2(n-2)}}\left(\int_{\Omega} u^{k(x)} d x\right)^{\frac{3}{4}}\left(\left(\int_{\Omega}\left|\nabla u^{\frac{k_{2}}{2}}\right|^{2} d x\right)^{\frac{n}{4(n-2)}}+\left(\int_{\Omega}\left|\nabla u^{\frac{k_{1}}{2}}\right|^{2} d x\right)^{\frac{n}{4(n-2)}}\right)
\end{gathered}
$$

Now, we can use the Young inequality to get

$$
\begin{gather*}
\int_{\Omega} u^{\frac{k(x)(2 n-3)}{2(n-2)}} d x \leq 2 B^{\frac{2 n}{3 n-8}} \frac{3 n-8}{4(n-2) \varepsilon^{\frac{n}{3 n-8}}}\left(\int_{\Omega} u^{k(x)} d x\right)^{\frac{3(n-2)}{3 n-8}}  \tag{2.16}\\
\quad+\frac{\varepsilon n}{4(n-2)}\left(\int_{\Omega}\left|\nabla u^{\frac{k_{1}}{2}}\right|^{2} d x+\int_{\Omega}\left|\nabla u^{\frac{k_{2}}{2}}\right|^{2} d x\right)
\end{gather*}
$$

where $\varepsilon$ is a positive constant to be determined later. Combining (2.16) with (2.13), we obtain

$$
\varphi^{\prime}(t) \leq C_{1}+C_{2} \varphi(t)^{\frac{3(n-2)}{3 n-8}}+C_{3}\left(\int_{\Omega}\left|\nabla u^{\frac{k_{2}}{2}}\right|^{2} d x+\int_{\Omega}\left|\nabla u^{\frac{k_{1}}{2}}\right|^{2} d x\right)
$$

where

$$
\begin{gathered}
C_{1}=\frac{1}{\left(k_{1}-1\right)^{2}}\left(\frac{1}{2} C_{k}+E(0)+\frac{1}{p_{1}+1} e^{\frac{p_{2}+1}{k_{1}-1}}|\Omega|\right)+\frac{|\Omega|}{k_{1}-1} \sup _{\Omega}\left(\alpha_{3}+\alpha_{1}\right) \\
C_{2}=\frac{4}{k_{1}-1}\left(\sup _{\Omega} \alpha_{2}\right) B^{\frac{2 n}{3 n-8}} \frac{3 n-8}{4(n-2) \varepsilon^{\frac{n}{3 n-8}}}, \\
C_{3}=\frac{2}{k_{1}-1} \frac{\varepsilon n}{4(n-2)}\left(\sup _{\Omega} \alpha_{2}\right)-\frac{2}{\left(k_{2}\right)^{2}}
\end{gathered}
$$

If we choose $\varepsilon>0$ such that

$$
0<\varepsilon \leq \frac{4(n-2)\left(k_{1}-1\right)}{\left(\sup _{\Omega} \alpha_{2}\right) n\left(k_{2}\right)^{2}}
$$

then, we obtain the differential inequality

$$
\begin{equation*}
\varphi^{\prime}(t) \leq C_{1}+C_{2} \varphi(t)^{\frac{3(n-2)}{3 n-8}} \tag{2.17}
\end{equation*}
$$

Integration of the differential inequality (2.17) from 0 to $t$ leads to

$$
\begin{equation*}
\int_{\varphi(0)}^{\varphi(t)} \frac{d \gamma}{C_{1}+C_{2} \gamma^{\frac{3(n-2)}{3 n-8}}} \leq t \tag{2.18}
\end{equation*}
$$

In fact, let $t \rightarrow T^{-}$, (2.18) leads to

$$
\int_{\varphi(0)}^{+\infty} \frac{d \gamma}{C_{1}+C_{2} \gamma^{\frac{3(n-2)}{3 n-8}}} \leq T
$$

where

$$
\varphi(0)=\int_{\Omega} \frac{1}{k(x)(k(x)-1)} u_{0}^{k(x)} d x
$$

Thus, the proof is achieved.
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