

Nonconvex second-order damped gradient systems and metastability

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find minimizers u^* of $E(\cdot)$

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- ▶ Gradient Flows $\frac{du}{dt} = -\nabla E(u)$,
- ▶ Second order Gradient Systems $\frac{d^2u}{dt^2} + \frac{du}{dt} = -\nabla E(u)$ or
- ▶ Systems with vanishing damping $\frac{d^2u}{dt^2} + \frac{\gamma}{t} \frac{du}{dt} = -\nabla E(u)$.

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AIM: discuss second order dynamical systems associated to a **nonconvex, noncoercive** functional E on an **infinite dimensional** space X .

A nonconvex functional in image processing

Samson et. al.¹ (cf. also Aubert and Kornprobst²) have proposed a Mumford-Shah-type nonconvex, **noncoercive** functional that can achieve both image classification and restoration simultaneously

$$E_\varepsilon(u) = \int_\Omega (u - u_0)^2 dx + \varepsilon \int_\Omega \varphi(|\nabla u|) dx + \frac{1}{\varepsilon} \int_\Omega W(u) dx$$

where u_0 is the image to be restored and classified, $\varepsilon > 0$ is a small parameter and $\varphi : \mathbb{R} \rightarrow \mathbb{R}$, $\varphi(w) = \frac{w^2}{w^2+1}$ while W is the double-well potential $W : \mathbb{R} \rightarrow \mathbb{R}$, $W(u) = \frac{1}{4}(u^2 - 1)^2$.

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Why not $\|u_x\|^2$?

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Why not $\|u_x\|^2$? Not edge-preserving (too much smoothing).

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Semilinear wave equations $u_{tt} + u_t = u_{xx} + f(u)$: classical results of Haraux & Jendoubi³

If f analytic then it satisfies the Lojasiewicz inequality.

Solutions satisfy the energy balance equation (with $F' = f$)

$$\frac{d}{dt} \left(\frac{1}{2} \|v\|^2 + \frac{1}{2} \|u_x\|^2 + \int_{\Omega} F(u) dx \right) = - \|v\|^2.$$

Theorem (Haraux-Jendoubi, '01)

Assume that $f : \mathbb{R} \rightarrow \mathbb{R}$ is analytic, with f'' is bounded on $(-\beta, \beta) \forall \beta > 0$ and let u be a solution such that

$$\bigcup_{t \geq 1} [u(t), v(t)] \text{ rel. compact in } H^2(\Omega) \times H^1(\Omega)$$

then there exists a stationary point u^* $T > 0$ large enough such that for all $t \geq T$

$$\|u(t) - u^*\|_{H^1} \leq Ct^{\theta/(1-2\theta)} \quad \text{if } 0 < \theta < \frac{1}{2}$$

$$\|u(t) - u^*\|_{H^1} \leq Ce^{-\omega t} \quad \text{if } \theta = \frac{1}{2}.$$

³A. Haraux, M. Jendoubi, *Decay estimates for some evolution equations with an analytic nonlinearity*, *Asympt. Anal.* 26(2001), 21-36.

A semilinear wave equation for the image processing functional

The semilinear equation associated to

$$E(u) = \int_{\Omega} \varphi(u_x) dx + \int_{\Omega} W(u) dx$$

is

$$u_{tt} + u_t = \varphi''(u_x)u_{xx} + u - u^3,$$

with the energy-dissipation equation

$$\frac{d}{dt} \left(\frac{1}{2} \|v\|^2 + \int_{\Omega} \frac{u_x^2}{1 + u_x^2} dx + \int_{\Omega} W(u) dx \right) = - \|v\|^2.$$

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NO... even well-posedness fails. Discontinuities develop in finite time.

Possible remedies/Alternative approaches

- ▶ weakening the solution concept (?!?)

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- ▶ modifying the energy functional by adding some higher order regularization (???)
- ▶ changing the damping (more/stronger damping)
- ▶ both changing the damping and adding a higher order regularization to the functional

Approach I: stronger damping + h.o. regularization

We consider

$$E(u)_\varepsilon = \frac{\varepsilon}{2} \|u_{xx}\|^2 + \int_{\Omega} \varphi(u_x) dx + \int_{\Omega} W(u) dx$$

and the higher order, damped wave equation

$$u_{tt} = u_{ttxx} - u_{xxxx} + \varphi''(u_x)u_{xx} + u - u^3.$$

Energy balance now reads

$$\frac{d}{dt} \left(\frac{1}{2} \|v\|^2 + \frac{\varepsilon}{2} \|u_{xx}\|^2 + \int_{\Omega} \frac{u_x^2}{1+u_x^2} dx + \int_{\Omega} W(u) dx \right) = -\|v_x\|^2.$$

Approach I (continued): rigorous results based on semigroup methods

This approach has a series of advantages:

- ▶ the linear part of the equation generates an analytic, immediately compact semigroup (see Engel & Nagel)
- ▶ global solutions exist in $H_0^2(\Omega) \times L^2(\Omega)$
- ▶ with improved regularity $(u(t), v(t)) \in H^4(\Omega) \cap H_0^2(\Omega) \times H_0^2(\Omega)$ for all $t > 0$
- ▶ all trajectories are relatively compact (by a result of Pazy)
- ▶ all trajectories converge to equilibrium (by LaSalle's Invariance Principle)
and furthermore
- ▶ the model admits an ε -independent $\|u_{xx}\|$ -estimate, on bounded time intervals $[0, T]$, such that solutions converge (as $\varepsilon \rightarrow 0$) to weak solutions of the $\varepsilon = 0$ -model (by Aubin's Lemma)

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but also one drawback

- ▶ the $t \rightarrow \infty$ and $\varepsilon \rightarrow 0$ limits do not commute.

Approach II: strong damping

We consider the unperturbed energy

$$E(u) = \int_{\Omega} \varphi(u_x) dx + \int_{\Omega} W(u) dx$$

and the damped wave equation with a strong damping

$$u_{tt} = u_{txx} + \varphi''(u_x)u_{xx} + u - u^3.$$

Energy balance now reads

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Approach II (continued): rigorous results

This analysis of this model is much more involved than what we had previously:

- ▶ local solutions exist in $H_0^2(\Omega) \times L^2(\Omega)$ (by semigroup methods or by the fixed point theorem of Krasnoselskii for the sum of two operators⁴)
- ▶ global existence in $H_0^2(\Omega) \times L^2(\Omega)$ follows from a new $\|u_{xx}\|$ -estimate
however
- ▶ no improved regularity or compactness is present
- ▶ one can not apply LaSalle's Invariance Principle.

⁴G. Andrews, *On the existence of solutions to the equation*
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Work in progress: energy decay estimates along bounded trajectories.

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Returning to the classical semilinear wave equation: Metastability and the role of small coefficients

Let us return to the classical semilinear wave equation, but with a small coefficient $\varepsilon \ll 1$

$$u_{tt} = -u_t + \varepsilon^2 u_{xx} + u - u^3.$$

There exists a manifold \mathcal{M} of initial data, in the state space, for which theorem of Haraux and Jendoubi still holds but with a very large T

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$$T = T_\varepsilon \geq \frac{C}{\sqrt{\varepsilon}} e^{1/\varepsilon}.$$

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This phenomenon is called Metastability⁵

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